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OPEN Temperature-dependent iron motion in extremophile rubredoxins – no need for 'corresponding states'

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Extremophile organisms are known that can metabolize at temperatures down to - 25 °C (psychrophiles) and up to 122 °C (hyperthermophiles). Understanding viability under extreme conditions is relevant for human health, biotechnological applications, and our search for life elsewhere in the universe. Information about the stability and dynamics of proteins under environmental extremes is an important factor in this regard. Here we compare the dynamics of small Fe-S proteins - rubredoxins - from psychrophilic and hyperthermophilic microorganisms, using three different nuclear techniques as well as molecular dynamics calculations to quantify motion at the Fe site. The theory of 'corresponding states' posits that homologous proteins from different extremophiles have comparable flexibilities at the optimum growth temperatures of their respective organisms. Although 'corresponding states' would predict greater flexibility for rubredoxins that operate at low temperatures, we find that from 4 to 300 K, the dynamics of the Fe sites in these homologous proteins are essentially equivalent.

Keywords Rubredoxin, Iron-Sulfur, Extremophile, Hyperthermophile, Psychrophile, Corresponding States

Over the past 4 decades, our knowledge about life under extreme conditions has dramatically expanded^{1,2}. Organisms have been found that can function down to -25 °C (psychrophiles) and up to 122 °C (hyperthermophiles)³. Microbial activity has been demonstrated at GPa pressures⁴, and growth at 100 MPa can has also been demonstrated (piezophiles)^{5,6}, as well as over a wide range of pH values (acidophiles and alkaliphiles)⁷, ionic strengths (halophiles)⁸, and other stressors and even combinations thereof⁹. These living systems must modulate their protein properties to function best in their preferred environment. Understanding how these extremophiles do so is interesting in its own right. It is also relevant for human health, biotechnological applications, and our search for life elsewhere in the universe¹⁰.

The dynamics of proteins in their proteins is an important part of extremophile adaptation. The temperaturedependent flexibility of enzymes is relevant to the optimum growth temperature for an organism and the range over which it can survive. Here the flexibility is defined as the rms deviation of atomic positions: $\sqrt{\langle \mu^2 \rangle}$. According to the 'corresponding states' hypothesis, enzymes have evolved to have comparable flexibility at the ideal growth temperature for their respective organism^{11,12}. In a similar vein, it is often stated that "directed thermal motion is needed for catalysis" and that "flexibility is necessary to allow catalysis at a metabolically appropriate rate"^{13,14}. In support of these ideas, neutron scattering experiments found comparable flexibility for psychrophiles and thermophiles at their respective adaptation temperatures (from 4 to 85 °C)¹⁵.

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However, the 'corresponding states' model is not universally accepted, and some experiments conflict with its predictions. For hyperthermophilic *Pyrococcus furiosus* rubredoxin (*Pf* Rd), NMR-monitored amide hydrogen exchange experiments found larger flexibility comparable to the mesophile *Clostridium pasteurianum* (*Cp*) protein^{16,17}. Using molecular dynamics calculations, Grottesi and coworkers found that at 300 or 373 K, hyperthermophilic *Pf* Rd is more flexible than the homologous mesophilic Rd from *Clostridium pasteurianum*¹⁸. Yet, a decade later, Rader compared the same pair with MD and explained the greater thermostability of *Pf* Rd *vs. Cp* Rd as the result of a 'decrease in flexibility'¹⁹. In a similar vein, from neutron scattering, NMR, and other measurements, hyperthermophilic P450 CYP119 was found to be more flexible than its mesophilic counterpart CYP101A at all temperatures above 200 K²⁰. In view of the discordance between the 'corresponding states' theory and multiple experiments, we decided to investigate the flexibility of different Rds at the Fe site by a variety of methods.

Rubredoxins (Rds) are the simplest Fe-S proteins²¹; they have an Fe(S-Cys)₄ center and molecular masses on the order of 6 kDa (Fig. 1). Rubredoxins are electron transfer proteins²², and they play important roles in photosystem II assembly²³, as redox partners of alkane hydroxylases^{24,25} and P-450 enzymes²⁶, and as platforms for artificial catalysts^{27,28}. As such they are an ideal test system for studies of extremophile dynamics. The structures and dynamical properties of rubredoxins have been studied by a variety of techniques and over ~ 10 orders of magnitude time scales. The experimental methods vary from inelastic neutron scattering on the picosecond scale^{16,29} to amide H–D exchange experiments over tens of seconds^{30,31}. Rubredoxins have also been extensively investigated by molecular dynamics (MD)^{18,32–37}.

For this study we have compared Rds from the hyperthermophile *Pyrococcus furiosus* (*Pf*)³⁸ with those from the psychrotolerant organisms *Pseudomonas* sp. strain AU10 (*Px*)³⁹ and *Polaromonas glacialis* (*Pg*). *Pf* is a marine archaeon with an optimal growth temperature of ~ 100 °C⁴⁰. *Px* is an Antarctic freshwater organism with an optimal growth temperature of 28 °C that can still grow close to 0 °C³⁹. *Pg* has been isolated from alpine glaciers and grows well from 1 to 25 °C^{41,42}. For *Pf*, it is presumed that *Pf* Rd is reduced by NADPH rubredoxin oxidoreductase (NROR), and that *Pf* Rd in turn reduces either a superoxide reductase (SOR) or a peroxide-reducing rubrerythrin (Rr)⁴³. For the psychrophilic Rds, since these species are implicated in the biodegradation of hydrocarbons and xenobiotics, one can assume that *Pg* and *Px* Rd provide electrons to enzymes that metabolize such molecules.

We have quantified the amount of ⁵⁷Fe motion in these rubredoxins by measuring the Lamb-Mössbauer factor as a function of temperature. The Lamb-Mössbauer factor, f_{LM} , is the ratio of elastic line intensity to overall nuclear absorption. In the case of harmonic motion $f_{LM} = \exp(-k^2 \langle \mu^2 \rangle)$, where $k = 2\pi/\lambda$ and $\langle \mu^2 \rangle$ is the mean square motion of the resonant nucleus. This breaks down when the motion is anharmonic, but f_{LM} remains a useful metric.

We have derived f_{LM} from three different measurements as well as from molecular dynamics calculations. One approach has been Nuclear Resonance Vibrational Spectroscopy (NRVS), where the elastic and inelastic intensities can be directly compared^{44–47}. We also conducted conventional Mössbauer experiments on these Rds. In this case we obtained relative f_{LM} values over a range of temperatures and then calibrated an absolute value at the low temperature limit with respect to the NRVS results. A third approach, also using synchrotron radiation, obtains relative f_{LM} values from the intensity of the Nuclear Forward Scattering (NFS)^{48,49}. Finally, the values from these three experimental methods have been compared with each other and with molecular dynamics calculations. In preliminary feasibility studies, we found that the oxidized Fe(III) form of rubredoxin was susceptible to photoreduction at higher temperatures. Thus, for all of these initial dynamics studies, we have focused on the reduced Fe(II) form of rubredoxin, which did not show any spectral changes during data collection.

The original Mössbauer approach to protein dynamics⁵⁰ and its interpretation as a protein dynamical transition have been criticized^{51,52}. NRVS and NFS are alternative and complementary probes, and f_{LM} can be reliably extracted from both NRVS and NFS data⁵³. All three of these methods sense motion of a single isotopic type of nucleus, in this case Fe-57, compared to neutron scattering which averages over (mostly protons) for an entire protein or organism.

Experimental results

Nuclear resonance vibrational spectroscopy (NRVS)

Since the procedure for deriving Lamb-Mössbauer factors from NRVS is not that common in bioinorganic chemistry, we first illustrate the temperature-dependent ⁵⁷Fe PVDOS (Partial Vibrational Density Of States) and normalization procedures in Fig. 2. The data processing involves estimation and removal of the elastic peak



Figure 1. (A) Ball and stick picture of the $Fe(S-Cys)_4$ site in Pf Rd. (B) Pymol diagram accentuating aromatic amino acids in Pf Rd. (C) Ribbon diagrams of Pf Rd with homology structure of Px Rd.



Figure 2. (**A**) Raw NRVS at 37K for *Pf* Rd before subtraction of elastic peak and normalization (Red line), scaled so that vibrational structure is visible. The (Black line) curve is the 37K data reduced by a factor of 100. (**B**) Processed NRVS data after removal of the elastic peak and normalization to PVDOS: reduced *Pf* Rd 46K (Red line), reduced *Pg* Rd 48K (Blue line), oxidized *Pf* Rd 46K (Red dotted line), oxidized *Pg* Rd 52K (Blue dotted line).

due to conventional Mössbauer scattering and a Fourier transform procedure that removes multiphonon events and yields a normalized PVDOS⁵⁴.

In Fig. 3A,B, we compare NRVS-derived values for the Lamb-Mössbauer factors and $\langle u^2 \rangle$ for the ⁵⁷Fe sites in all three rubredoxins over a range of temperatures. The trend for *Pf* Rd is smooth and approximately linear over most of the temperature range. There is no evidence for a so-called 'dynamical transition' involving a rapid change in slope in the region above 200 K, as derived from conventional Mössbauer spectroscopy for myoglobin^{53,55,56}. Finally, within experimental error, at comparable temperatures, the $\langle u^2 \rangle$ is the same for the psychrotolerant *Px* Rd, the psychrophilic *Pg* Rd, and the hyperthermophilic *Pf* Rd.

In Fig. 3C, the NRVS-derived $\langle u^2 \rangle$ for *Pf* Rd is compared with theoretical curves for Debye models⁵⁷ with a variety of temperatures. There is good agreement with a Debye temperature of ~175 K. Previous work by Debrunner and coworkers put the Rd Debye temperature between 145 and 212 K⁵⁸. Using 175 K, the Debye model predicts $\langle u^2 \rangle = 0.00375$ Å² at 4.2 K and hence a low-temperature Lamb-Mössbauer factor $f_{LM} = 0.82$. This compares favorably with the Lamb-Mössbauer factor $f_{LM} = 0.85 \pm 0.06$ calculated for *Cp* rubredoxin from 4.2 K Mössbauer data by Debrunner and coworkers⁵⁸. The temperature dependence of their data was also consistent with our results.

The calculated mean square Fe motion in Fig. 3B is predicated on the assumption of harmonic motion for the ⁵⁷Fe nucleus, and anharmonic effects can require corrections^{59–62}.

Temperature-dependent Rubredoxin Mössbauer

One difference between the NRVS-derived $\langle u^2 \rangle$ and the previous Mössbauer results for Rd⁵⁸ or heme proteins^{50,63} is the time scale of the experiments. NRVS reflects a time scale on the order of 4 picoseconds, while the lifetime of a Mössbauer measurement is on the order of 100 ns⁶³. Accordingly, we also ran the conventional Mössbauer for reduced *Pf* Rd and *Pg* Rd, and the results are presented in Fig. 4. At 4 K the Rd spectra are virtually identical (Fig. 4A) with respective quadrupole splittings of 3.28–3.31 mm s⁻¹ and equal isomer shifts of 0.67 mm s⁻¹. With increasing temperature, the absorption as governed by f_{LM} decreases dramatically (Fig. 4B,C). It has often been



Figure 3. (A) NRVS-derived Lamb-Mössbauer factors *vs*. T for ⁵⁷Fe in *Pf* Rd (Red filled circle), *Pg* Rd (Blue filled circle), and *Px* Rd (Green filled diamond). (B) Comparison of NRVS-derived $< u^2 >$ for the ⁵⁷Fe sites in *Pf* Rd (Red filled circle), *Pg* Rd (Blue filled circle), and *Px* Rd (Green filled diamond). Straight line is linear fit to just *Pf* Rd data. (C) $< u^2 > vs$. temperature for different Debye temperatures as labeled, compared with experimental values from *Pf* Rd NRVS (Red filled circle).



Figure 4. (A) Low temperature (4.2 K) Mössbauer spectra for reduced Pf (Red line) and Pg (Blue line) Rd. The Pg Rd intensity was scaled to match Pf Rd, but the velocity scales were unchanged. (B) Representative Mössbauer spectra for Pf Rd at different temperatures. (C) Lamb-Mössbauer factors f_{LM} for Pf (Red filled circle) and Pg (Blue filled circle) Rd *vs.* temperature, using f_{LM} normalized to 0.82 at 4.2 K. Pf and Pg data superimpose at 4 K and 180 K.

noted, that, "precise measurement of f_{LM} is difficult with conventional Mossbauer spectroscopy⁴⁸", and "it is difficult to determine absolute values for the recoilless fractions ... without tedious and difficult experimentation⁶⁴." So, the scatter in the observed values is not unexpected.

The Mössbauer data for both rubredoxins shows a dramatic falloff around 235 K. In this regard it is consistent with previous work on Mb that posited a dynamical transition above 200 K^{55,56,63}. Thus, as with the previous Mb work, we are seeing significantly different mean square Fe motion on the time scales of the Mössbauer and NRVS measurements. We note that the time-scale of the Mössbauer measurement is governed by the ~ 100 ns excited state lifetime. In contrast, it has been argued that the NRVS time scale is on the order of 4 ps⁶³.

To obtain an absolute value for the Lamb-Mössbauer factor, we used the low temperature (4.2 K) limit of the Debye model with a Debye temperature of 175 Kelvin (Fig. 3C). This gave $\mu^2 = 3.66 \times 10^{-3} \text{ Å}^2$ and hence $f_{LM} = 0.82$. This value is reasonably close to the $4.9 \times 10^{-3} \text{ Å}^2$ value for $< u^2 >$ derived from the lowest temperature NRVS data at 37 K. We thus feel justified in using $f_{LM} = 0.82$ as the low temperature limit with which to calibrate the 4 K Mössbauer data. Our 0.82 value is similar to the $f_{LM} = 0.85 \pm 0.06$ reported by Debrunner et al. for oxidized *Pf* Rd as far back as 1979⁵⁸. This allows us to report absolute f_{LM} in Table 1.

Temperature-dependent Rubredoxin nuclear forward scattering

Nuclear forward scattering is a synchrotron technique that provides Mössbauer-like information from the timedependent scattering of the emitted radiation. Interference between the different scattered frequencies produced by quadrupole and/or magnetic interactions generates 'quantum beats' in the scattered intensity, and there can also be 'dynamical beats' in the scattered intensity due to absorption and reemission of the beam as it passes through the sample⁶⁵. As early as 1994, Bergmann and coworkers used NFS to extract temperature dependent f_{LM} for Fe metal⁴⁸.

	Linewidth (mm -s ⁻¹)		$\Delta E_Q (mm s^{-1})$				Lamb-Mössbauer Factor*			
T (Kelvin)			<i>Pf</i> Rd		Pg Rd		<i>Pf</i> Rd		Pg Rd	
	<i>Pf</i> Rd	Pg Rd	Möss	NFS	Möss	NFS	Möss	NFS	Möss	NFS
4.2	0.28	0.28	3.28		3.31		0.82		0.82	
18				3.24		3.27		0.82		0.84
40	0.28	0.27	3.28	3.24	3.32	3.27	0.78	0.77	0.69	0.78
90	0.28	0.27	3.26	3.23	3.30	3.26	0.63	0.65	0.59	0.66
100				3.22		3.26	0.62	0.63	0.58	0.63
120	0.27	0.26	3.25	3.21	3.28	3.25	0.60	0.58	0.55	0.58
150	0.27	0.26	3.24	3.20	3.27	3.24	0.59	0.51	0.40	0.51
180	0.27	0.26	3.23	3.19	3.27	3.23	0.53	0.43	0.38	0.43
200	0.27	0.26	3.22	3.18	3.26	3.22	0.44	0.39	0.50	0.38
220	0.26	0.26	3.21	3.17	3.26	3.21	0.42	0.34	0.32	0.33
235	0.40	0.40	3.20	3.17	3.24	3.20	0.17	0.25	0.34	0.25
250								0.04	≤0.05	0.04

Table 1. Mössbauer and NFS Results on Reduced Fe(II) *Pf* and *Pg* Rd \sim *Assuming $\mu^2 = 3.6 \times 10^{-3}$ Å² at 4 K. \sim Interpolated as necessary.

The properties of the scattered radiation depend on the effective sample thickness χ , where $\chi = \eta_s d \sigma_0 f_{LM}$. Here, η_s is the volume density of elastic scatterers, d is the sample thickness, and σ_0 is the maximum resonance cross section⁶⁵. In the case of a quadrupole split Mössbauer spectrum, there are two scattered frequencies with splitting $\Delta \omega$, and the transmitted intensity $I_{tr}(t)$ can be described as a function of time t by⁶⁵:

$$I_{tr}(t) \propto \exp(-\tau) \frac{\chi}{\tau} J_1^2 \left(\sqrt{0.5 \chi \tau} \right) \cos^2 \left(\frac{\Delta \omega}{2} t + \frac{\chi \Gamma_0}{8 \Delta E} \right)$$
(1)

In this equation, τ is the time in units of the nuclear lifetime τ_0 and J_1 is a Bessel function of the first kind. The quadrupole splitting leads to a beat time⁶⁵:

$$T_{qb} = \frac{2\pi}{\Delta\omega} = \frac{86ns}{\Delta E_Q[mms^{-1}]}$$
(2)

Hence, the quadrupole splitting can be obtained from the beats: $\Delta E_0 [\text{mm s}^{-1}] = 86 \text{ ns}/T_{\text{qb}}$.

The NFS signals for both Rds are compared in Fig. 5A,B. A least-squares fit of the *Pg* Rd NFS data at 18 K to a version of Eq. (1) is shown in Fig. 5C. The optimized value $\Delta \omega = 0.240$ yields a $\Delta E_Q = 3.27$ mm s⁻¹, compared with the low temperature Mössbauer value of 3.31 s⁻¹. At the other temperature extreme, fitting the *Pf* Rd NFS data at 235 K yields $\Delta E_Q = 3.17$ mm s⁻¹ (Fig. 5D), compared to the Mössbauer value of 3.20 mm s⁻¹.

As seen in Eq. (1), the dependence of the NFS signal amplitude on f_{LM} is more complicated. As described in Supporting Information, the bottom line is that for concentrated samples over long times, the NFS signal under these conditions becomes approximately proportional to the Lamb-Mössbauer factor. However, for relatively thin samples, $\chi \leq 1$, and at short times, the NFS signal will be approximately proportional to χ^2 and hence to the square of the Lamb-Mössbauer factor: $I(t) \sim \chi^2 \exp[-(1 + \chi)\tau]^{66}$.

The temperature dependence of the NFS signals for both Rds is compared in Table 1, Fig. 5A,B,E,F. We see that there is a precipitous drop in the signal around 235 K, just as observed in the Mössbauer data. In other work the same effect has been seen in the NFS of myoglobin⁶⁷, and it has been taken as evidence for a 'dynamical transition'⁶³ or 'glass transition'^{51,68}. Of note is that this transition is not observed in our NRVS data. The agreement between NFS and Mössbauer measurements likely reflects that they explore the same ~ 100 ns time scale, compared to the much shorter time scale for NRVS.

Apart from f_{LM} , the NFS reveals the temperature dependence of the quadrupole splitting ΔE . Here we see the same trend as in the Mössbauer, but our values tend to be ~ 1% smaller than obtained by the conventional technique.

Molecular dynamics calculations for T-dependent ⁵⁷Fe motion

The amount of motion at the ⁵⁷Fe site can also be derived from molecular dynamics calculations. As computer technology has improved, such studies of rubredoxin have examined dynamics on progressively longer time scales: from 10 ps in 1993³², to 30 ps³⁵, 100 ps³⁴ 400 ps³⁶ 6 ns¹⁸ and recently out to 1 μ s Sala³⁷. Our current results on the 4 ns time scale are illustrated in Fig. 6.

Over the entire temperature range, from cryogenic temperatures to 300 K, there is relatively little difference in the calculated Fe motion between both proteins. This is in agreement with the experimental results from the previous nuclear spectroscopies. The amount of motion is slightly higher than obtained from NRVS, this may reflect the 4 ns timescale for the calculations *vs.* the ps timescale for NRVS. A brief look at 4 ps time-scale calculations found ⁵⁷Fe rmsd were reduced by ~ 20%, bringing them more in line with NRVS.

Summary of different approaches

The values for root mean square ⁵⁷Fe disorder derived from different techniques are summarized in Table 2. At 235 K, supposedly above the PDT temperature, there is excellent agreement between NRVS, Mössbauer, and MD techniques for reduced *Pf* Rd ⁵⁷Fe RMSD: 0.16 ± 0.01 Å. The agreement is reasonably good at lower temperatures. However, above 235 K, the signals from Mössbauer and NFS fall sharply, whereas the NRVS is still measurable.

In comparison, the RMSD values obtained via B-values from x-ray and neutron diffraction structures are much larger than the spectroscopic or MD determinations. For example, in the region near 100 K, all of the spectroscopic rmsd estimates are at or below 0.1 Å, whereas the smallest crystallographic values are larger: 0.20 Å for reduced *Pf* Rd at 100 K⁶⁹ and 0.17 Å for oxidized *Pf* Rd at 100 K⁶⁹. The pattern repeats at 295 K, where the diffraction values (0.33 Å or 0.34 Å) are again about twice the estimates from NRVS (0.17 Å) or MD calculations (0.20 Å). The larger diffraction values arise because these methods are sensitive to both static and dynamic disorder in the Fe positions.

Discussion

In this work we have compared two psychrophilic rubredoxins, Pg Rd and Px Rd, with a hyperthermophilic homologue from Pf Rd. In contrast with previous work which has primarily focused on the dynamics of protein protons, we have used nuclear spectroscopies that solely measure the dynamics of the labeled ⁵⁷Fe site. According to the 'corresponding states' hypothesis, the hyperthermophilic Rd should be less flexible than the psychrophilic protein at the same temperature.

Instead, we find that over a wide temperature range, from cryogenic temperatures up to ambient conditions, the dynamics at the Fe redox site are quite similar for both psychrophilic and hyperthermophilic proteins. Our results are thus consistent with hydrogen exchange studies by Hernández et al. that found no evidence "systematic rigidification" in *Pf* Rd¹⁶. In their MD simulations, Grottesi et al. even found greater flexibility for *Pf* Rd compared to mesophilic *Clostridium pasteurianum* (*Cp*) Rd¹⁸. In another HD exchange study comparing *Pf* and *Cp*



Figure 5. Nuclear forward scattering data. (**A**) Pg Rd (in order of decreasing intensity) measured at 18 K (Black line), 60 K (Red line), 110 K (Blue line), 170 K (Green line), and 235 K (Yellow line). Max overall count rate was ~ 1.4×10^7 s⁻¹. Bin-width = 0.4675 ns. (**B**) Same for Pf Rd. (**C**) Least squares fit (Red line) and data (Black line) for Pg Rd NFS at 18 K. (**D**) Least squares fit (Red line) and data (Black line) for Pf Rd NFS at 235 K. (**E**) Normalized count rate for Pf (Red filled circle) *vs.* Pg (Blue filled circle) Rd. (**F**) Same on log scale.

Rds, LeMaster et al. found "no necessary correlation between thermostabilization and increased conformational rigidity"⁷⁰. These results along with our findings differ from the MD calculations of Rader who found a decreased flexibility for Pf Rd compared to Cp Rd⁷¹.

Of course, our use of nuclear probes that focus on motion at the Fe site raises question about the broader implications for protein dynamics of the whole protein. Although the Fe is situated at one end of the molecule, in previous work we have shown that Fe participates in low frequency modes that "*involve concerted in-phase collective motion of large segments of polypeptide*^{72,73}." Others have proposed that even the Fe–S modes are "*extensively coupled to deformations of the polypeptide backbone*^{74,75}." Thus, although our probes were Fe-centered, rubredoxins are such small proteins that Fe motion should be a good proxy for global protein dynamics.

Our results raise the question: how important is flexibility for the electron transfer role of rubredoxin? Although there are as yet no structures for our Rds bound to redox partners, we can learn from the *Pseudomonas aeruginosa* system. The *Pa* rubredoxin, Rdx, receives electrons from an FADH-based reductase, RdxR, and



Figure 6. Molecular dynamics predictions for ⁵⁷Fe motion as a function of temperature. Time scales for these calculations were 4 ns. (**A**) Average of 3 MD runs for Pf Rd (Red filled circle) *vs.* 1 run for Pg Rd (Blue filled circle). Smooth curve is quadratic fit to Pf Rd calculations to guide the eye and allow interpolation for Table 2.

T T (Kelvin)	Diffraction	NRVS	Mössbauer	NFS	molecular dynamics [§]
4	-	0.06 [‡]	0.06	-	-
40	-	0.075	0.068	0.067	0.068
100	0.17 ^{5x} 0.20 ^{† x}	0.10	0.074	0.093	0.10
110	0.23 ^{& x}	0.10	0.10		0.11
123	0.28 ^{Ω x}	0.11	0.10		0.11
160	0.30 ^{z x}	0.12	0.12	0.12	0.13
200	0.27 ^{z x}	0.13	0.13	0.14	0.15
235	0.34 ^{z x}	0.14	0.18	0.16	0.17
250		0.15	-	-	0.18
295		0.16	-	-	0.20

Table 2. Reduced Fe(II) *Pf* Rubredoxin T-Dependent ⁵⁷Fe RMS Disorder (Å) by Various Techniques. ^xxray diffraction. ⁿneutron diffraction ^zunpublished *Pf* Rd [‡]normalized to Debye curve. [†]PDB 5OME [∞]PDB 4AR6. [§]PDB 4AR4 [‡] **PDB 2DSX PDB 3KYU ^Ω 1BRF [&]2DSX (*Dg* Rd). [‡]Extrapolated from Debye model. [§]Interpolated from smooth curve fitted to average of 3 simulations.

transfers them to alkane hydroxylases AlkB1 and AlkB2²⁴. The structure of a complex between Rd and RdxR has been solved to 2.45 Å by x-ray crystallography⁷⁶. Using an AlfaFold model for the isolated Rdx, we estimate that the average structural change in the Rdx molecule before and after binding RdxR is on the order of 0.5 Å⁷⁶. The flexibility associated with typical enzymatic conformational changes does not seem required for Rd to accomplish its mission. Indeed, Hageleuken et al. have observed "the docking of Rdx precisely at the tunneling hot spot of RdxR, and the absence of any appreciable conformational changes during Rdx binding"⁷⁶.

Flexibility is frequently reported as key for the functioning of enzymes. As expressed by Tsou, "flexibility ... is mandatory for the maximal expression of enzyme activity"⁷⁷. Or, as articulated by Hammes and coworkers, flexibility is one of the "Pillars of Enzyme Catalysis"⁷⁸. However, it is not clear how much flexibility is required for electron transfer proteins such as rubredoxin. Rubredoxin needs to dock with its respective electron donors and acceptors, but it may be that the electrostatic forces, perhaps combined with partner flexibility, are sufficient to bring the Rd Fe center close enough for electron transfer.

Evidence for Rd promiscuity can be seen in *Mycobacterium tuberculosis*, where the native rubredoxin (RubB) can shuttle electrons from two cognate reductases, FprA and FdR, to several different heme-based P-450 enzymes: CYP124, CYP125, and CYP142. Sushko et al. concluded that the rubredoxin interactions were "transient and not highly specific"⁷⁹.

To reiterate, much of the discussion about the importance of appropriate flexibility of motion *vs.* temperature, relates to conformational changes needed to enable enzyme activity. However, we note that the rubredoxins that are the subject of the current study are electron transfer agents and conformational changes may not be critical for their activity. Thus, if the natural partners of the current rubredoxins merely require proximity of the Fe site to the donor or acceptor site, flexibility might not be a key determinant for its effectiveness. Perhaps we should consider rubredoxins as closer to reagents than to enzymes.

We note that the optimal growth conditions for organisms, their protein flexibility, and the thermal stability of their constituent proteins, while correlated, are mathematically related. Data for melting temperatures of psychrophile rubredoxins are scant, so instead we refer to the well-studied mesophile *Clostridium pasteurianum* rubredoxin – *Cp* Rd. Here the optimized growth temperature for *Cp* at pH 6.5 is 40 °C⁸⁰, while T_m value for *Cp* Rd was reported as 104 °C³. This contrasts with the optimal growth temperature of 100 °C for *Pyrococcus furiosus*, while the T_m value for *Pf* Rd was reported as 104 °C³

Of course, flexibility might still be important for the range of temperatures over which Rd is stable, but that will involve thermodynamic issues instead of kinetic ones. Although the thermal stability of Pg Rd as not been investigated, it is known that mesophile Cp melts around 57 °C⁸¹, while Pf Rd melts at 144 °C⁸². NRVS at temperatures above 300 K will be difficult, and Mössbauer and NFS are out of the question. Furthermore, the studies have demonstrated that flexibility is variable across proteins³¹, and our current study has been limited to dynamics at the Fe site. We will need to employ other methods for answering such questions.

Summary

We began this study in part as a test for the corresponding states hypothesis^{83–85}. We found that psychrophilic rubredoxins and hyperthermophilic have comparable flexibility at the Fe site from 4 K to 300. It may be that flexibility is not essential to rubredoxin electron transfer, but additional studies at higher temperatures, perhaps by other techniques such as SRCD, SAXS, and NMR, are needed to examine the relationship between flexibility and thermal stability.

Methods

Production and purification of the recombinant rubredoxins

The recombinant *P. furiosus* rubredoxin was expressed and purified essentially as described³⁸. Recombinant *P. glacialis* rubredoxin was purified in a similar way. Based on the amino acid sequence (WP_196868757.1) the gene was codon optimized for expression in *E. coli*, synthesized and cloned into plasmid pET24a by Genscript to generate plasmid pPglRd. The recombinant vector containing the gene that codes for a rubredoxin from *Pseudomonas* sp. AU10 (GenBank accession number OP536816) was synthesized with codon optimization for *E. coli* in GenScript. In order to substitute the natural abundance Fe with ⁵⁷Fe, all proteins were expressed and purified as described above, with cultures that were supplemented with 50 μ M ⁵⁷Fe. Full experimental details are in the Supporting Information.

NRVS measurements

NRVS spectra were measured and recorded using published procedures⁸⁶ at either BL09XU, BL19LXU, or BL35XU. These beamlines employ two monochromators: first, a high heat load monochromator (HHLM) to produce ~ 1.0 eV resolution at the ⁵⁷Fe nuclear resonance (14.41 keV) and second, a high-resolution monochromator (HRM) to reduce the bandwidth to 0.8 meV (6.5 cm⁻¹)⁸⁷. The photons from both the nuclear fluorescence and the internal conversion K shell fluorescence were recorded with a 2×2 APD array.

The ⁵⁷Fe partial vibrational density of states (PVDOS), the real temperatures (T_r), the Lamb-Mössbauer factors (LM) and other sum rule quantities⁴⁶ for the samples were derived from the measured NRVS using the PHOENIX software package⁸⁸. Full experimental details are in the Supporting Information.

Mössbauer measurements

Mössbauer spectra were recorded by using a home-built spectrometer equipped with a Janis Research (Wilmington, MA) SuperVaritemp Dewar that allows measurements in the temperature range from 1.5 to 250 K and a constant-velocity transducer. The temperature was monitored by using a calibrated resistance type temperature sensor (Lake Shore Cernox CX-SD). The isomer shifts are quoted relative to α -Fe foil at 298 K. The velocity transducer was calibrated using a sodium nitroprusside standard.

Nuclear forward scattering (NFS)

NFS measurements were performed at BL19LXU at SPring-8 using the same high-resolution monochromator and cryostat as NRVS measurements. Photons from the high-resolution monochromator were transmitted through the sample at about 1 m distance onto a 7-element APD detector array a further 1 m downstream. NFS spectra were measured using the SPring-8 D-mode time structure, comprising a single synchrotron pulse of 13 ps with an interval of 684.3 ns. The APD detector was capable of counting at 14 kHz (2 kHz per array element) and the incident photons were attenuated to optimize to this detector count rate using aluminum filters before the sample. Two kinds of measurement were taken, time spectra at a given temperature, and total count rate between ~ 18–600 ns as a function of temperature.

NFS spectra were simulated using custom spectrum fitting software. Lamb-Mössbauer factors were estimated from the both the simulated spectra and total count rates as described in the Supporting Information.

Molecular dynamics calculations

Molecular dynamics simulations were carried out using GROMACS version 2021.5⁸⁹⁹⁰. Protein coordinates were obtained from a previous X-ray structural determination of the Rd from *P furiosus* (PDB: 1BRF)⁹¹. The simulations utilized the CHARMM27⁹² forcefield for protein interactions. Van der Waals interactions for the cofactor were obtained from previous work on rubredoxin, using the nonbonding parameters for the bridging sulfurs and iron face³³. The Fe-(Cys)4 interactions were adapted from empirical NRVS data⁷² that were previously simulated by VIBRATZ⁹³. Charges for the Fe-(Cys)4 site were obtained from Density Functional Theory

(DFT) calculations. The DFT geometry optimizations were performed with the program ORCA version 4.2^{94} using the BP86⁹⁵⁻⁹⁷ Customized with 10% exact Hartree–Fock exchange as functional and Ahlrich's triple-zeta basis set Def2-TZVPP⁹⁸ at the ORCA grid6 integration grid level and the CP(PPP)⁹⁹ basis set for Fe atoms at the grid7 integration grid level. The DFT calculation used implicit solvation (CPCM) with $\varepsilon = 4$ and was corrected for dispersion effects (D3)^{95,96}. Initial coordinates were taken from the XRD structures for Fe-(Cys)4 in 1BRF, then subsequently geometry optimized with the positions of all Ca carbons fixed to their starting positions¹⁰⁰. The charges for the Fe-(Cys)4 moiety were then extracted by the CHarges from Electrostatic Potentials using a Grid (CHELPG) method¹⁰¹ implemented in ORCA. Individual charges were scaled to a summed net charge of -2, and integrated into the molecular dynamics forcefield. More complete details are in the Supporting Information.

Data availability

The data that support this study are available from the corresponding authors upon request.

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Author contributions

S.P.C. wrote the main manuscript text and directed the research program F.E.J. and A.S. prepared the Pf and Pg rubredoxin samples. J.J.M. and S.C.-S. prepared the Px rubredoxin samples. T. D. prepared Fig. 1 and assisted in data interpretation. J.X. and Y.G. conducted the Mössbauer measurements and analysis. S.J.G., N.N., Y.Y conducted the NFS measurements. S.J.G. analyzed the NFS data and assisted in data interpretation. H.X. and M.Y.H. analyzed the NRVS data. H.M. K.T. N.N. L.L. Y.Y. and H.X. collected the NRVS data. L.B.G. performed the MD calculations. All authors reviewed the manuscript.

Competing interests

The authors declare no competing interests.

Additional information

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